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14. ABSTRACT Combustion instability in liquid rocket engines can have severe consequences including degraded performance, accelerated component wear, and potentially catastrophic failure. High-frequency instabilities, which are generally the most harmful in liquid rocket engines, can be driven by interactions between disturbances associated with transverse acoustic resonances and the combustion process. The combustion response to acoustic perturbation is a critical component of the instability mechanism, and is in general not well understood. The current paper describes an experimental facility at the Air Force Research Laboratory (AFRL) at Edwards Air Force Base that is intended to investigate the coupling between transverse acoustic resonances and single/multiple liquid rocket engine injector flames. Critical aspects of the facility will be described, including the capability to operate at supercritical pressures that are relevant to high-performance liquid rocket engines, accurately-controlled and cryogenically-conditioned propellants, and optical access to facilitate the use of advanced diagnostics. The transverse acoustic resonance is induced through the use of carefully-controlled piezo-sirens, allowing monochromatic excitation across a range of amplitudes at a number of discrete frequencies. The location of the flame within the acoustic resonance mode shape can also be varied through relative phase control of the two acoustic sources. The operating space of the facility, for oxygen and hydrogen operation, will be described. Preliminary non-reacting and reacting data will also be presented to demonstrate the quality of operation of this facility. It is anticipated that future results generated using this facility will provide both fundamental insight into the acoustic-flame interactions as well as provide a database useful for validating combustion instability models.				
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Development of a Facility for Combustion Stability Experiments at Supercritical Pressure

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Supported by
Air Force Office Scientific Research
Air Force Research Laboratory

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Overview

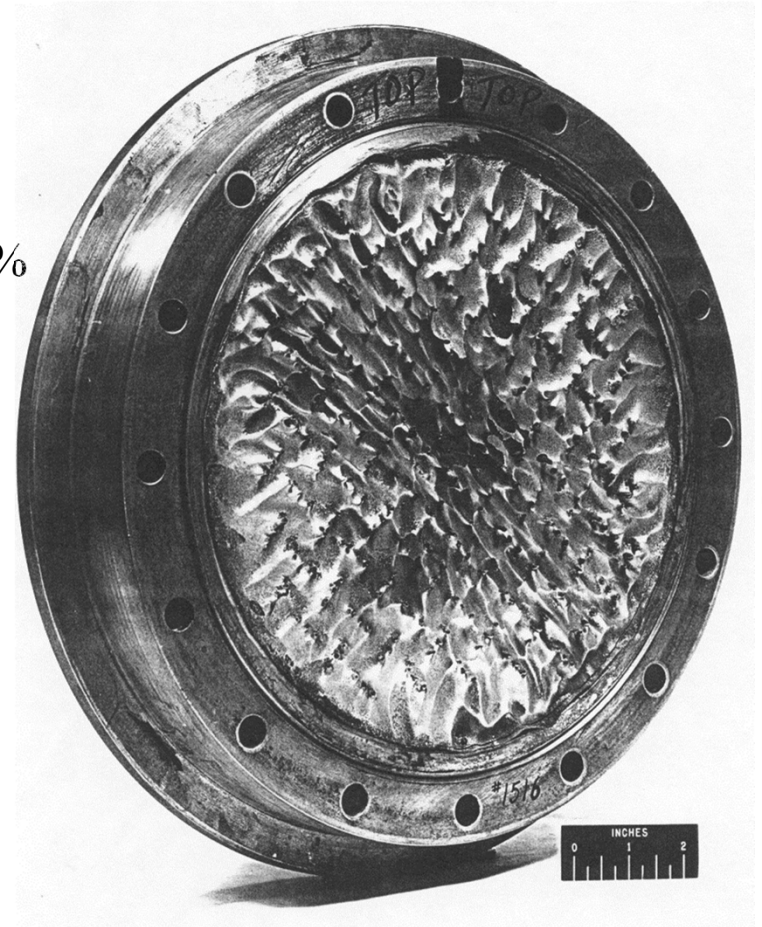
- Background
 - Combustion Instability
 - Challenges
- Experimental Techniques
 - Facility
 - Heat Exchangers
 - Acoustic Excitation System
 - Injector
 - Proper Orthogonal Decomposition of High-speed Images
- Preliminary Results

Background

- Liquid Rocket Engine Combustion Instability

Motivation

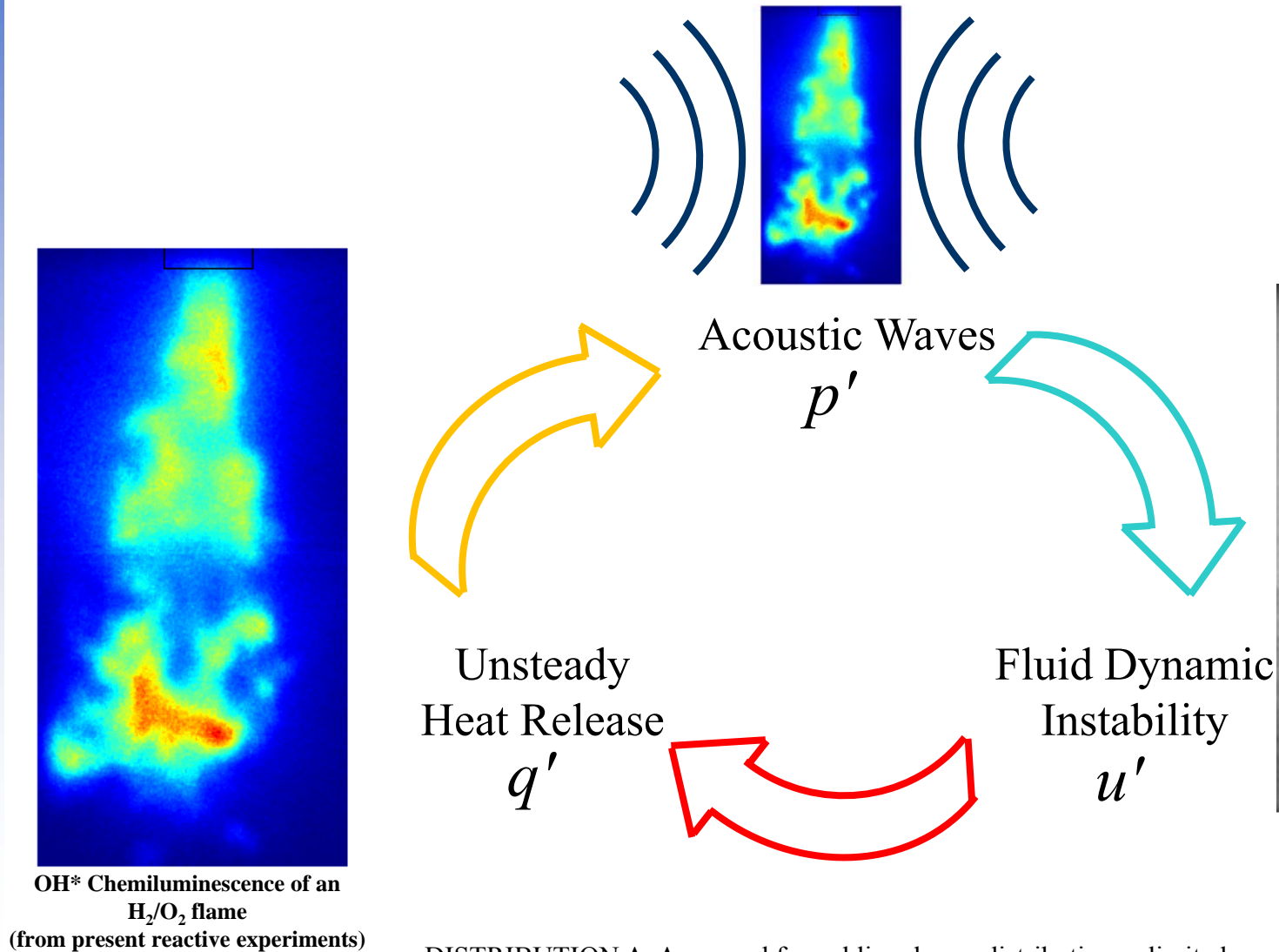
- An organized, oscillatory motion in the combustion chamber sustained by combustion
- Chamber pressure amplitudes p' can exceed 100% of the mean chamber pressure p_c
- Most difficult instabilities to eliminate:
high frequency (a.k.a. “screaming” instability)



Irreparable damage can occur in $< 1s$

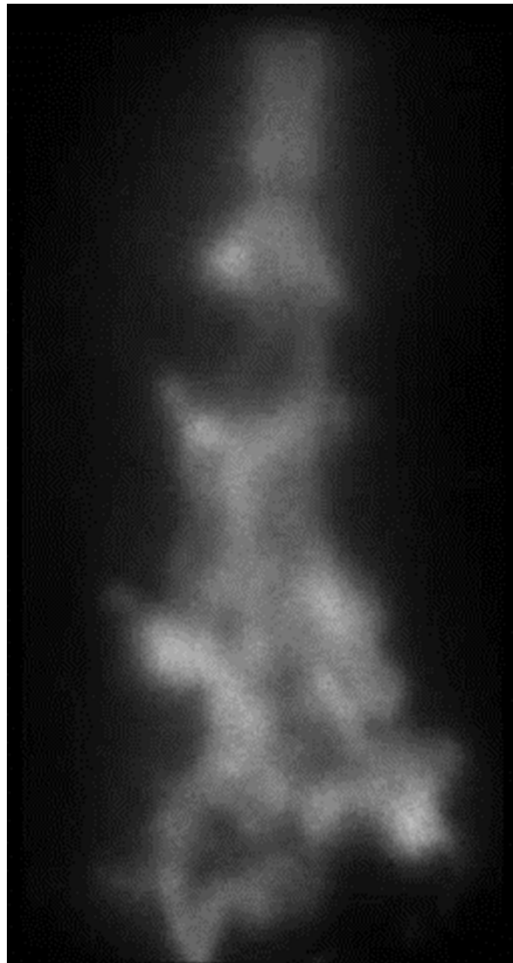
Background

- Combustion Instability Feedback Loop



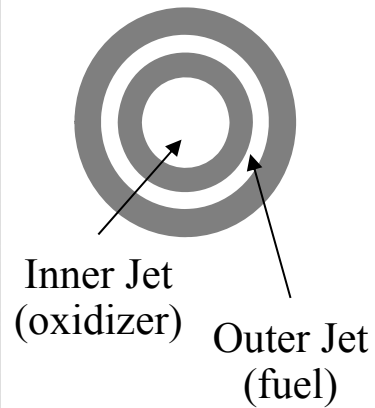
Background

Reactive



OH* Chemiluminescence
Coaxial O₂-H₂ Flame
 $p_c = 400$ psia

Coaxial Injector
Cross-sectional View



Density Ratio

$$S = \rho_{oj}/\rho_{ij}$$

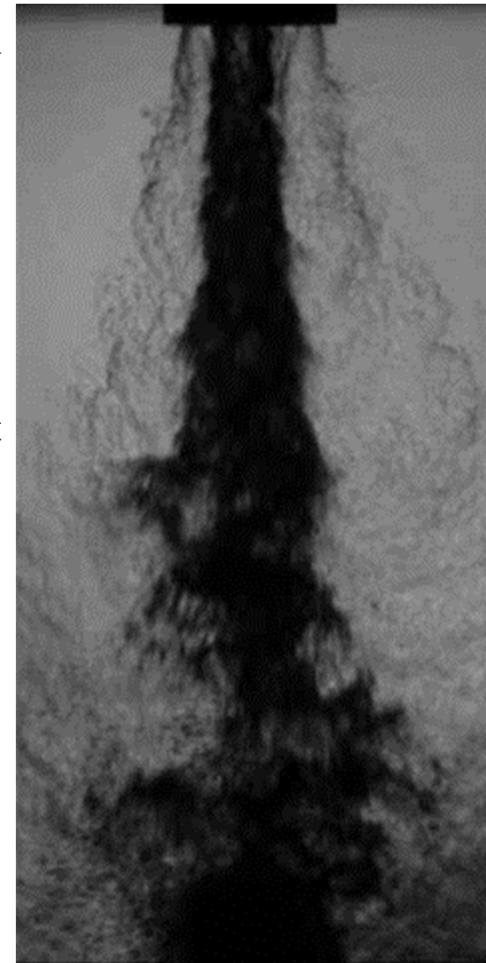
Velocity Ratio

$$R = u_{oj}/u_{ij}$$

*Momentum Flux
Ratio*

$$J = SR^2$$

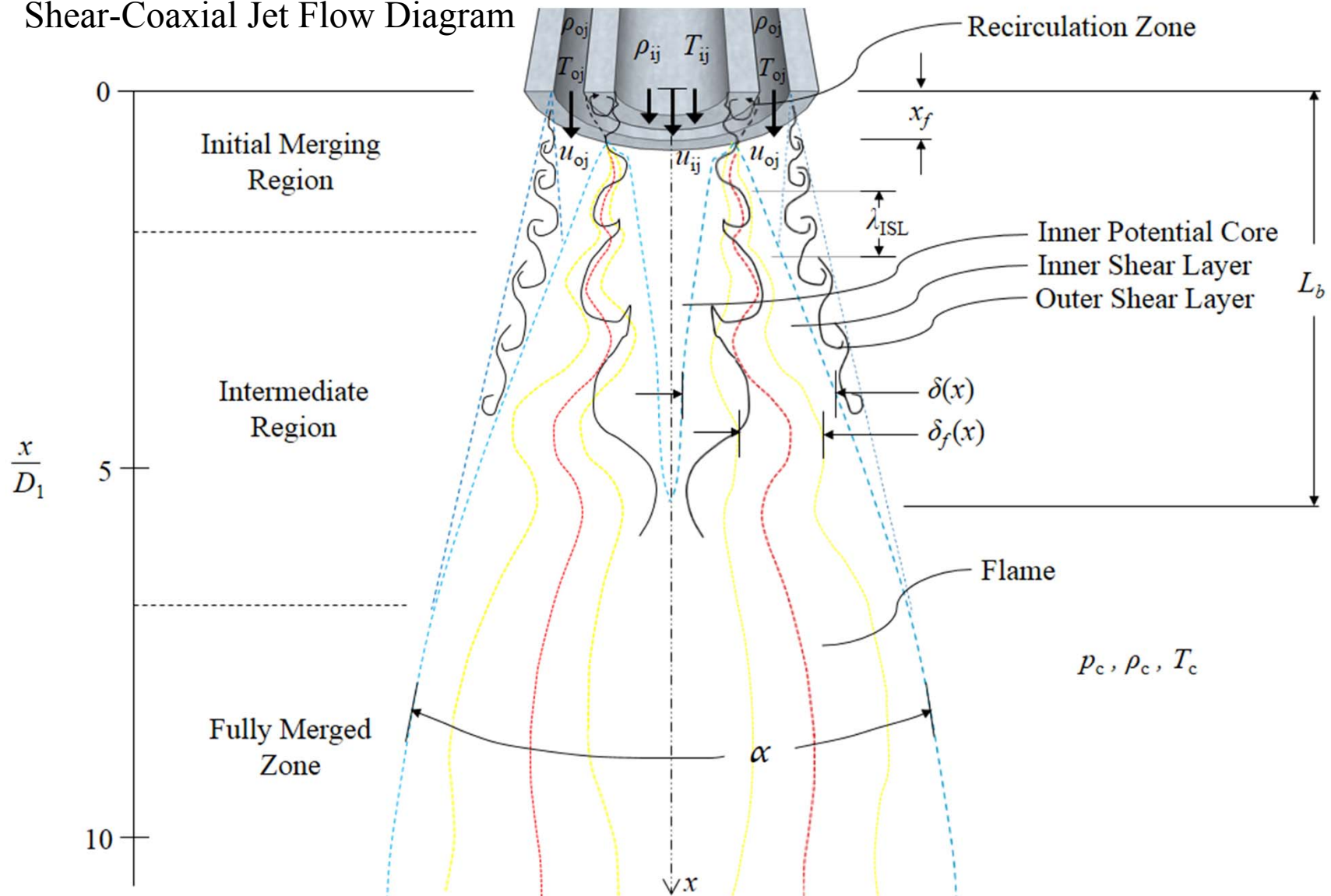
Nonreactive



Back-lit Imaging
Coaxial LN₂-GHe Jet
 $p_c = 400$ psia

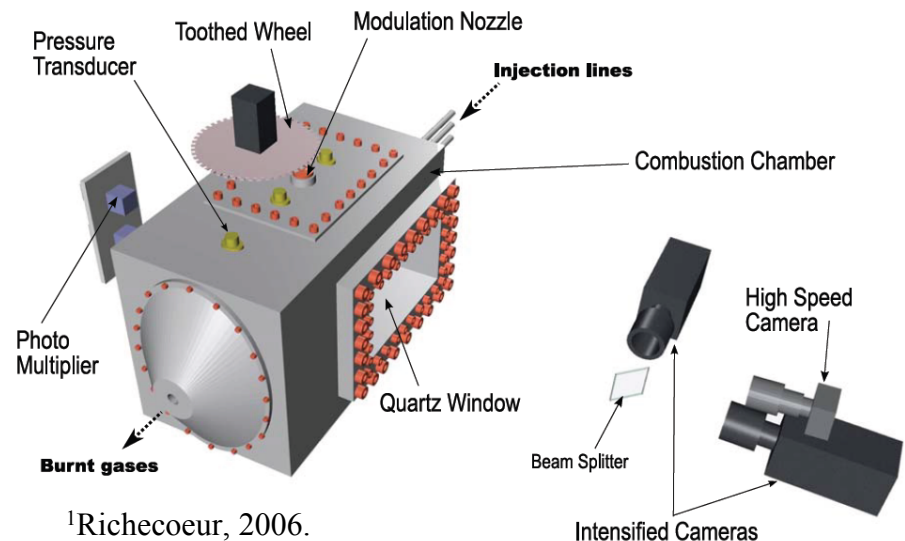
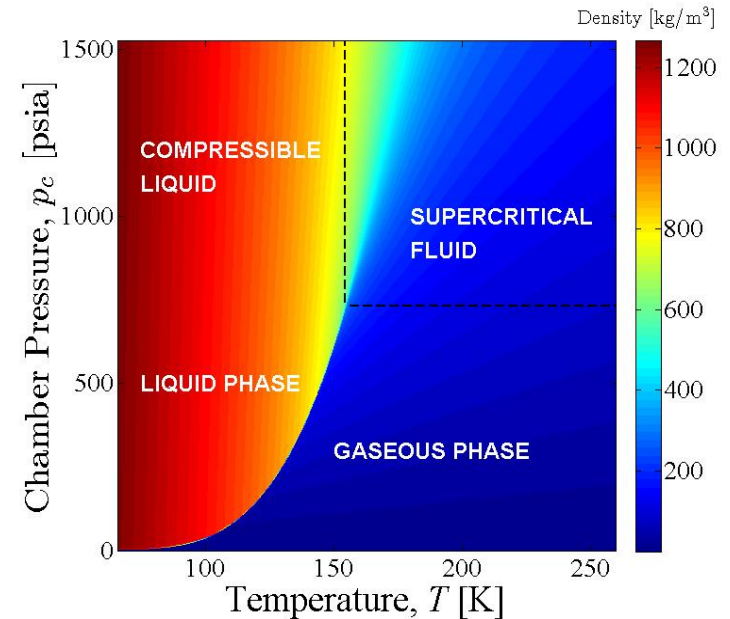
Background

- Shear-Coaxial Jet Flow Diagram



Background: Challenges Associated with Combustion Stability Experimental Facilities

- Supercritical Chamber Pressure
 - Liquid rocket engines (LREs) often employ chamber pressures greater than the critical pressure of oxygen
 $p_c > 731 \text{ psia}$
 - Surface tension and phase changes are undefined
- High Amplitude Acoustic Perturbations
 - Severe LRE combustion instabilities involve pressure amplitudes far greater than those of traditional acoustic excitation systems ($p' \sim p_c$)
- Imaging Diagnostics
 - High-speed ($>10 \text{ kHz}$) optical equipment is required to capture unsteady heat release via chemiluminescence and planar laser induced fluorescence (PLIF)
 - Windows must endure large changes in temperature as well as high chamber pressures



¹Richecoeur, 2006.

Existing Facilities

DLR (Germany)

CNRS (France)

Purdue Univ.

Penn. State Univ.

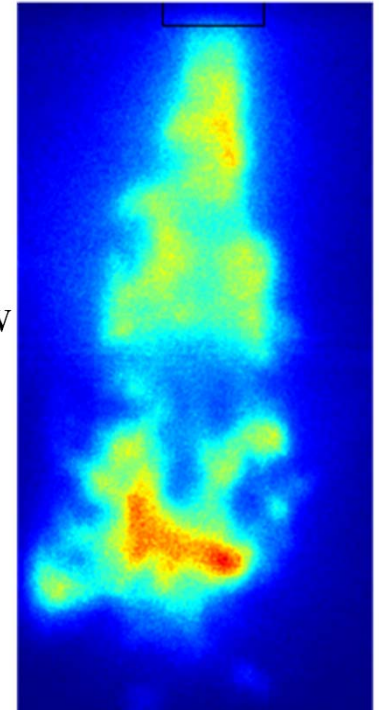
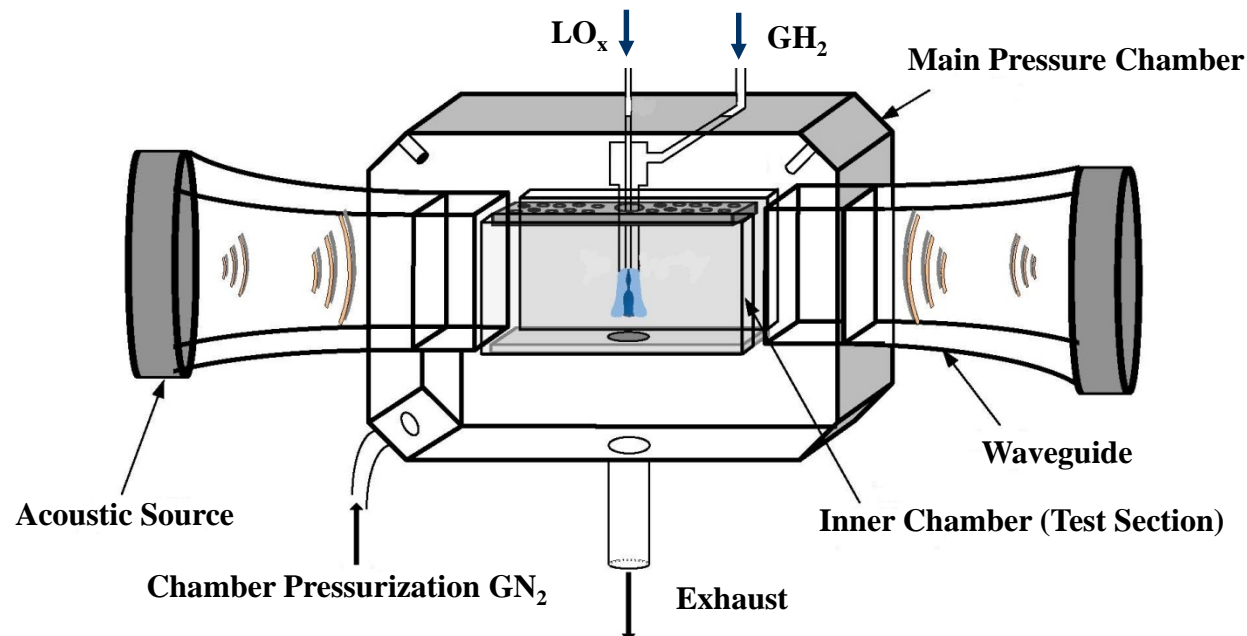
Experimental Techniques: Facility

- Capabilities

- Cryogenic propellant temperature control with high accuracy (± 1 K)
- Sub- and super-critical chamber pressure (p_c up to 10.4 MPa)
- High amplitude acoustic forcing ($p'/p_c \sim 0.02$)
- Coaxial injector with extended length for fully developed turbulent flow
- High speed diagnostic tools

Pressure transducer(s) natural frequency > 100 kHz

Time-series OH* chemiluminescence imaging ($f > 10$ kHz)



OH* Chemiluminescence
of an H₂/O₂ flame
(from present reactive
experiments)

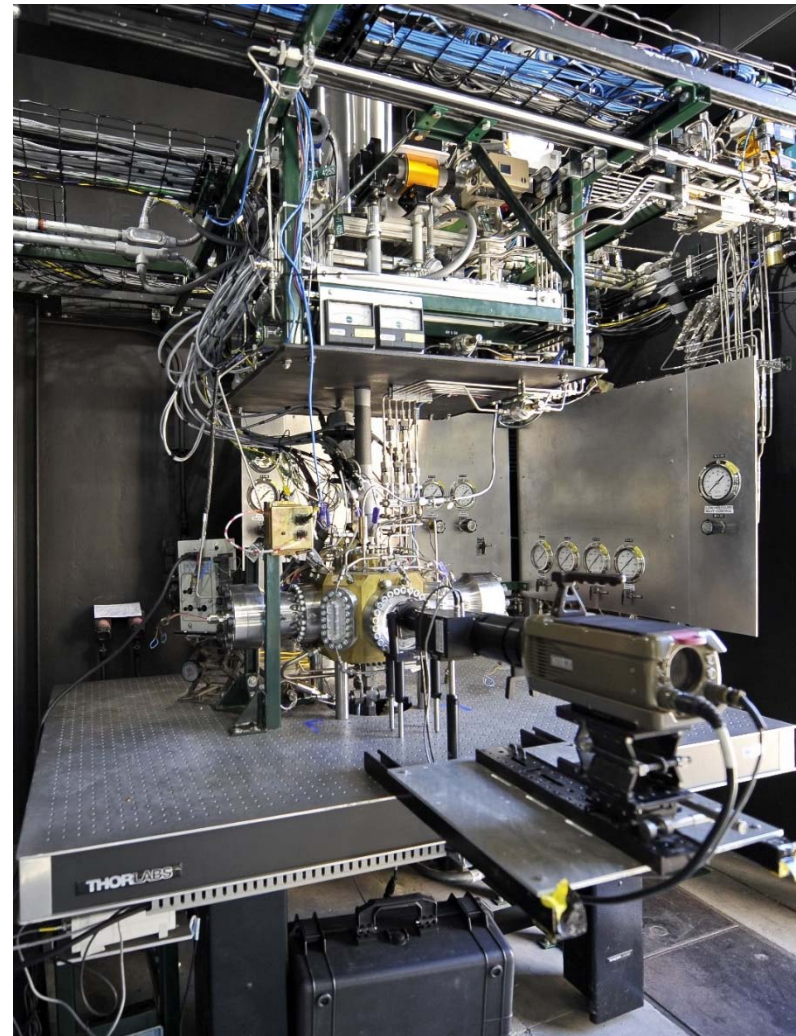
Experimental Techniques: Facility

BEFORE



2011

AFTER



2013

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Experimental Techniques: Facility

BEFORE



2011



AFTER



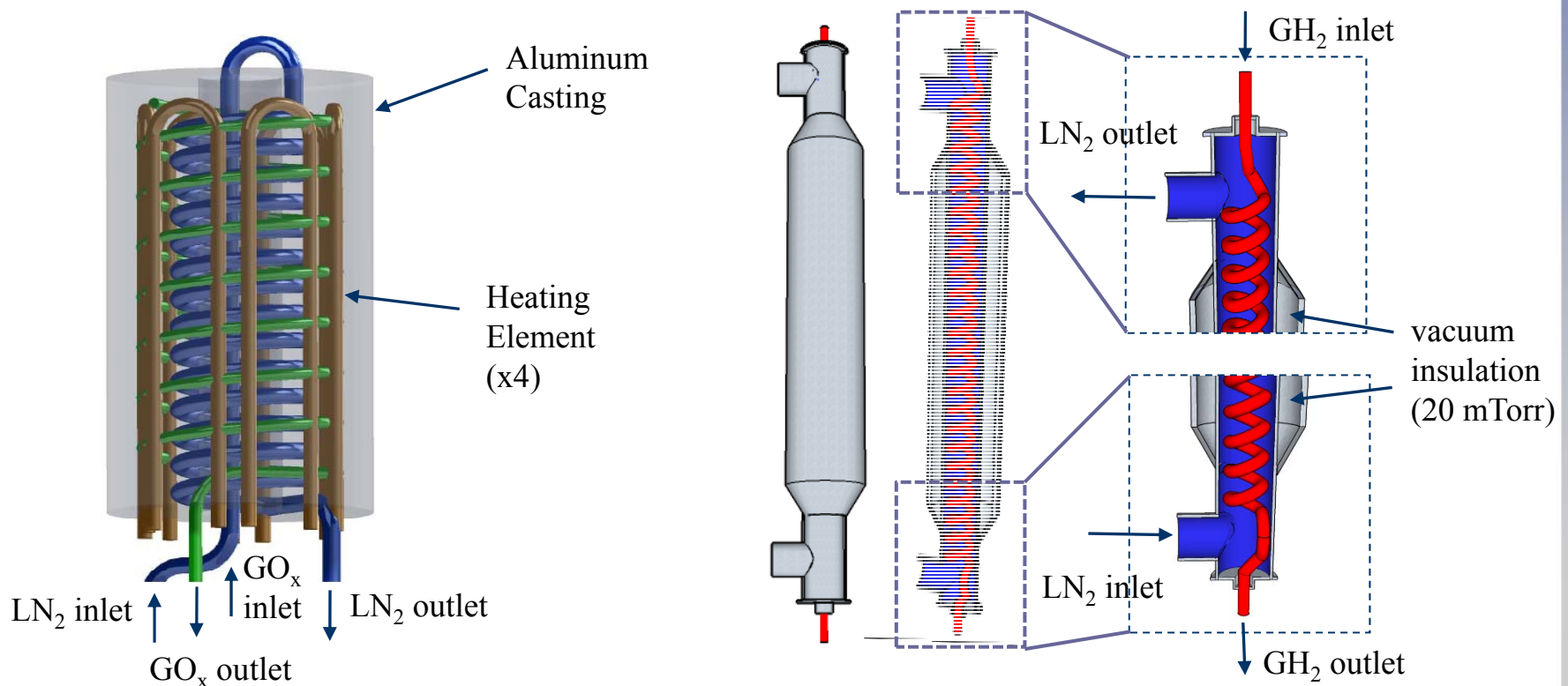
2013



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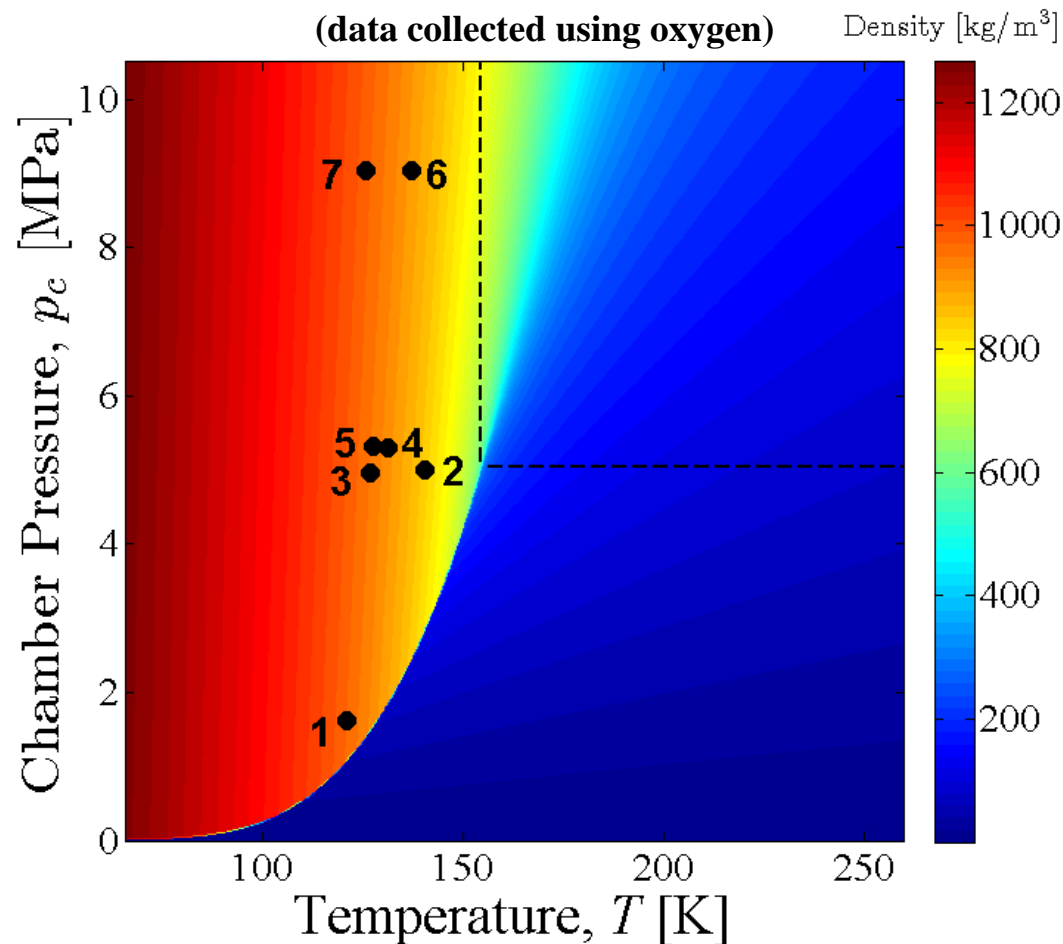
Experimental Techniques: Heat Exchangers

- Heat Exchanger Design Objectives
 - Achieve inner and outer jet temperatures, T_{ij} and T_{oj} , which are similar to the propellant temperatures of a liquid rocket engine
 - Minimize: temperature control error, coolant flow rate, user interaction



Experimental Techniques: Heat Exchangers

- Inner Jet Heat Exchanger



	\dot{m}_{ij}
• 1	3.7 g/s
• 2	2.6 g/s
• 3	9.3 g/s
• 4	3.0 g/s
• 5	11.1 g/s
• 6	2.7 g/s
• 7	10.4 g/s

Both gaseous oxygen and liquid oxygen conditions can be achieved.

Inner jet mass flow rate has a significant effect on inner jet temperature.

Pressure has a negligible effect on inner jet temperature.

Experimental Techniques: Waveguides

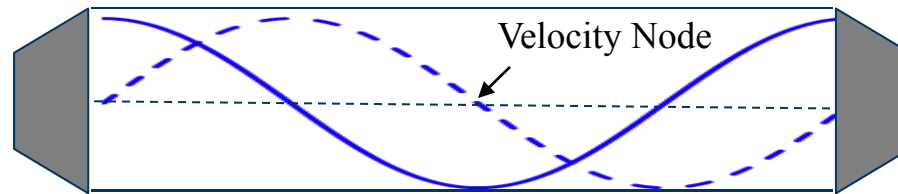
- Objective: To create an acoustic field with a **transverse standing wave**.

Velocity Node: $L = n\lambda = n \frac{c}{f_n}$

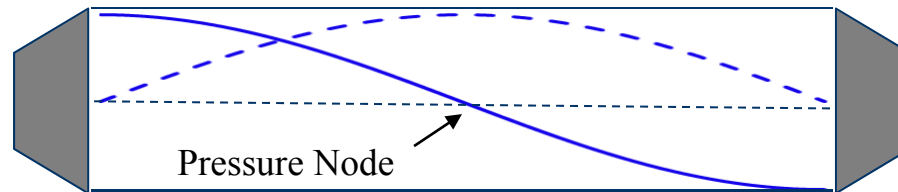
Pressure Node: $L = \left(\frac{2n+1}{2}\right)\lambda = \left(\frac{2n+1}{2}\right) \frac{c}{f_n}$

— p'
- - - u'

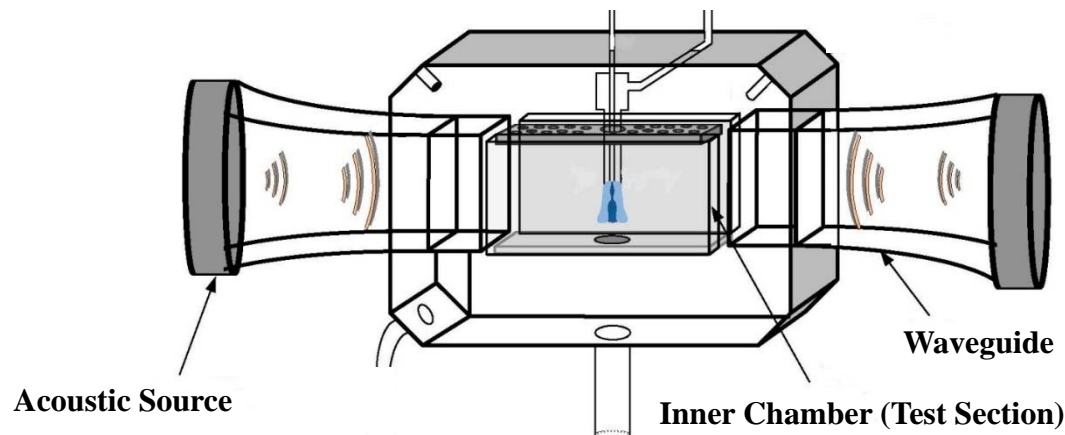
$$L = \frac{c}{2f_0}$$



$$\phi = 0^\circ$$



$$\phi = 180^\circ$$



Experimental Techniques: Waveguides

- Acoustic Waveguide Design

- Objectives

- Minimize: two- and three-dimensional waves
- Maximize: pressure amplitude

- Derivation of area relation

- Webster's horn equation

$$\frac{1}{A} \frac{\partial}{\partial y} \left(A \frac{\partial p}{\partial y} \right) - \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = 0 \quad ^{11}$$

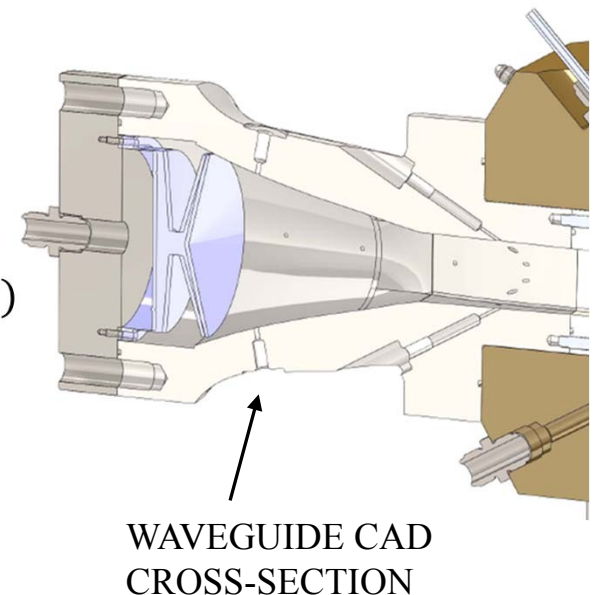
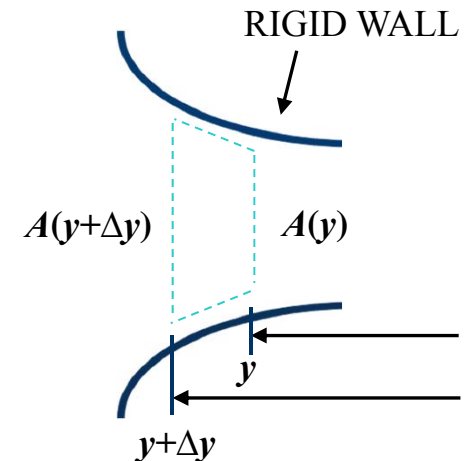
- General solution for area

$$A^{1/2} = A_{th}^{1/2} (\cosh my + T \sinh my)$$

where $A = A_{th} (\cosh my)^2$ and $A_{th}^{1/2} T m = \frac{\partial(A^{1/2})}{\partial y} (y = 0)$

- Particular solution for the catenoidal horn ($T = 0$)

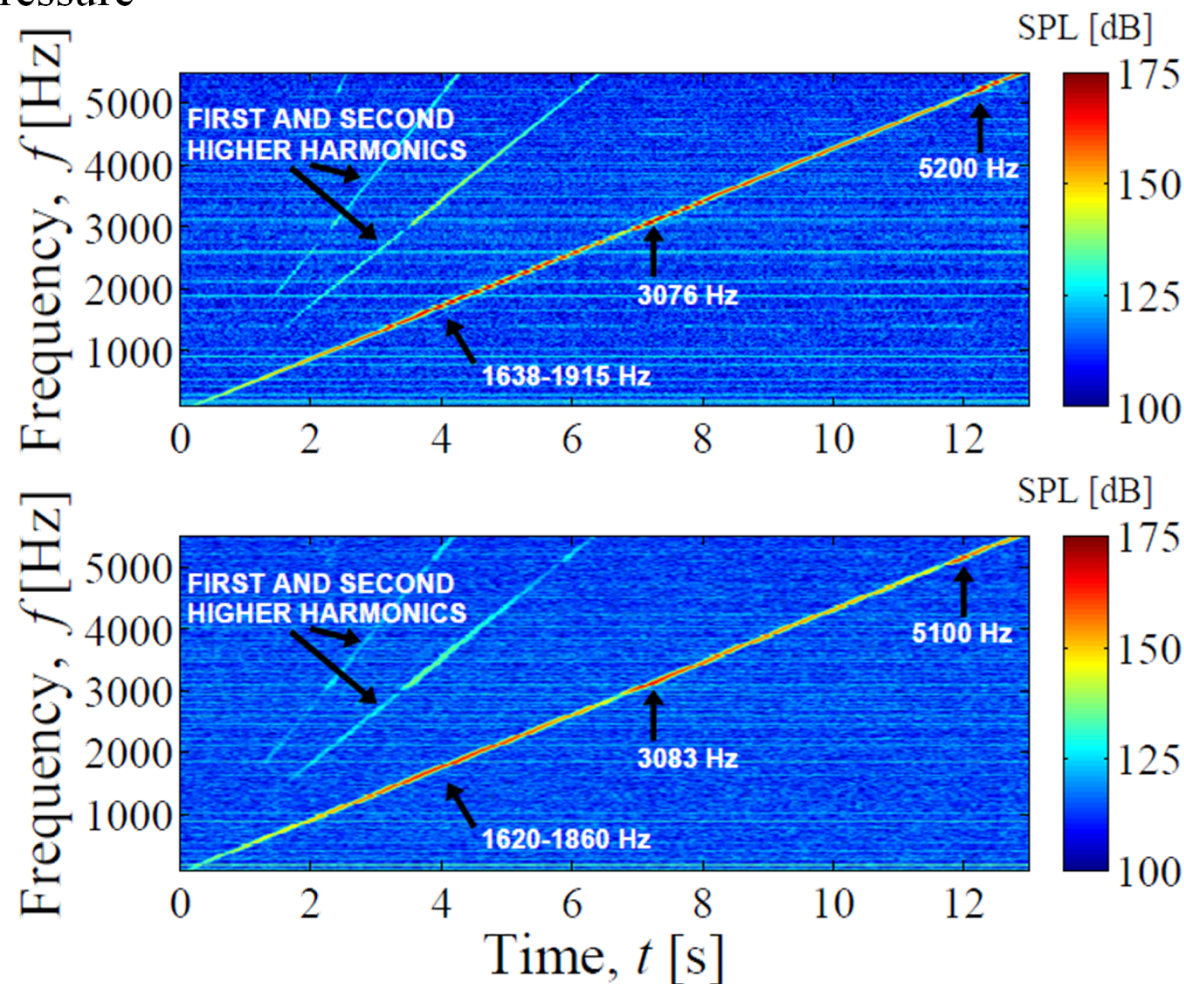
$$* \quad A = A_{th} (\cosh my)^2$$



² Pierce, A. D., Acoustics: An Introduction to Its Physical Principles and Applications, 2nd Edition, 360-363, 1991.

Experimental Techniques: Waveguides

- Acoustic Characterization: Piezoelectric sirens placed *outside* of the chamber at ambient pressure

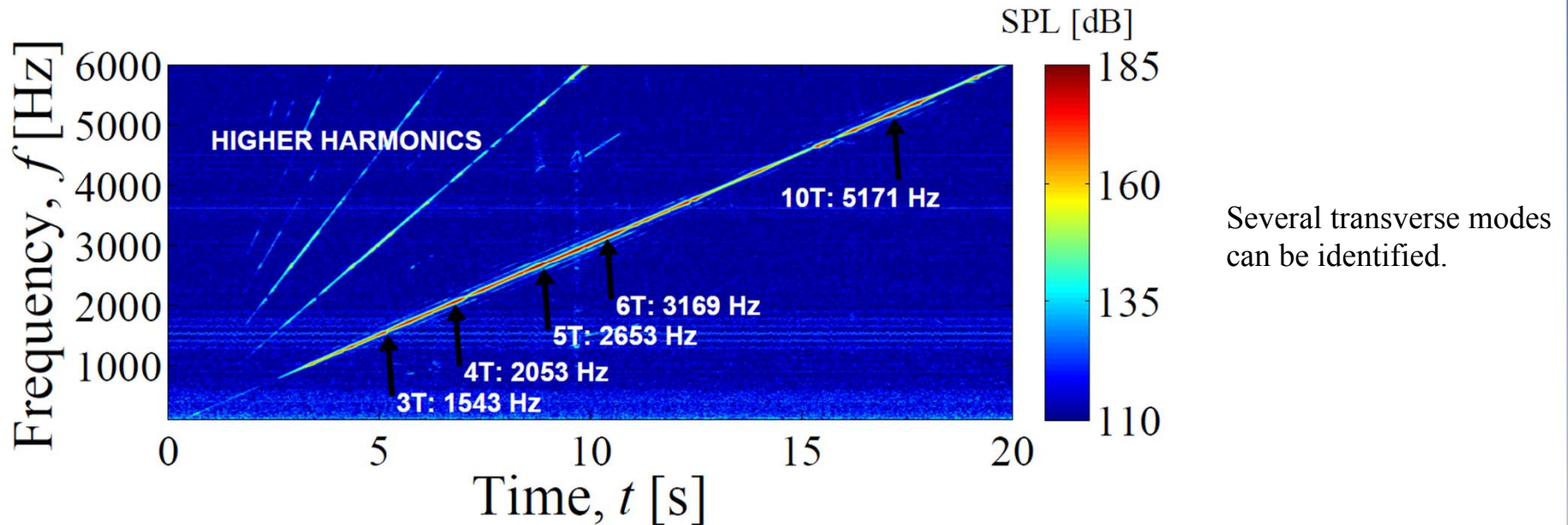


PIEZOELECTRIC SIREN

⇒ Each siren has a unique frequency response, with 3 peak frequency bands.

Experimental Techniques: Waveguides

- Acoustic Characterization: Piezoelectric sirens operated in-phase at $p_c = 400$ psia



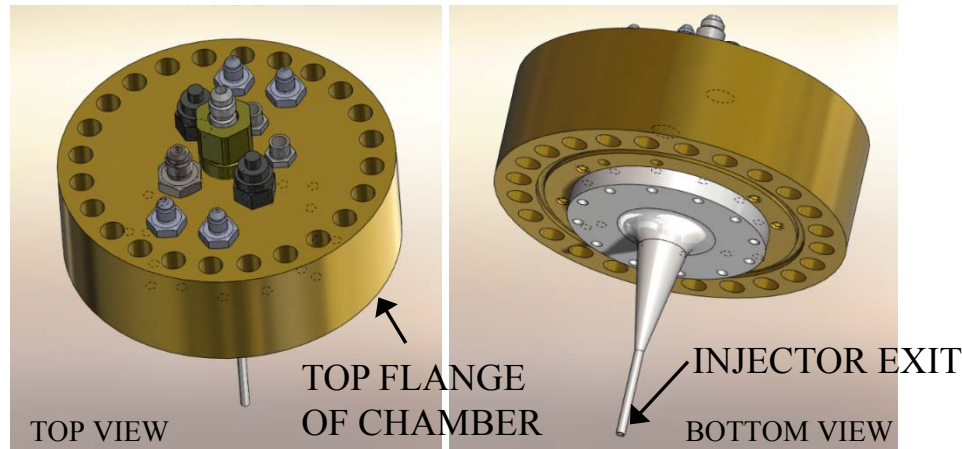
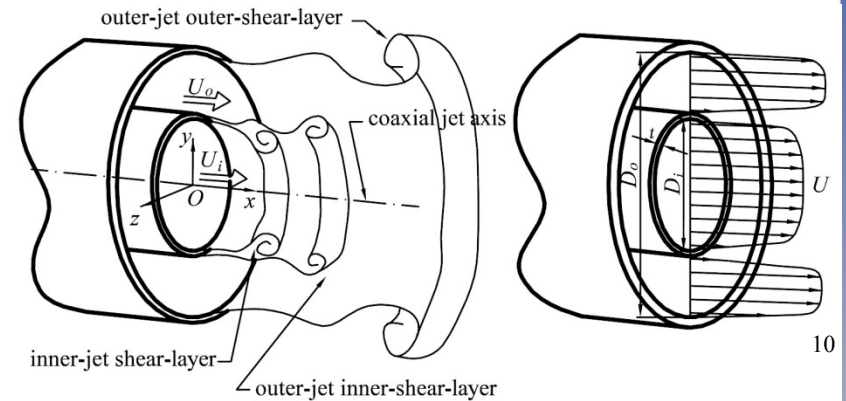
- Conclusions
 - An array of transverse modes will produce high amplitudes (3T, 4T, 5T, 6T, 10T modes).
 - Operating at other frequencies produces either low amplitudes, or undesired resonance, e.g. longitudinal resonance.
 - Peak frequency values are dependent on chamber pressure and temperature
- Standing waves must be verified prior to each individual experimental test.

Experimental Techniques: Injector

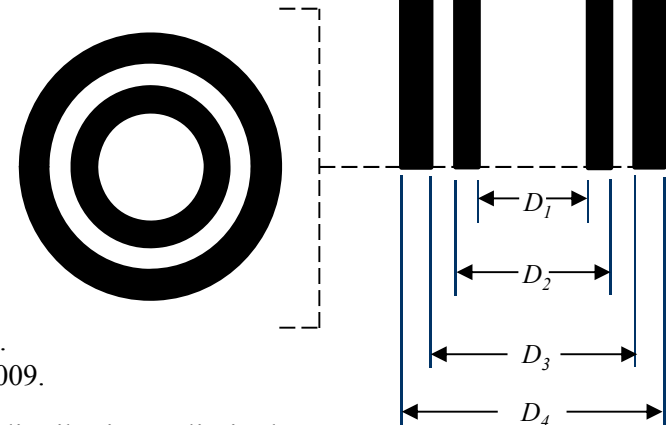
Objective: create a LO_x/GH_2 injector which provides fully turbulent exit flow

Challenges

- Materials
- Robust separation between GH_2 and LO_x
- Large injector length ($l_e/D \geq 4.4\text{Re}^{1/6}$ for fully developed turbulent flow⁹)
- Small cross sectional size
- Rigidity
- Manufacturability



$D_1 = 1.40 \text{ mm}$
 $D_2 = 2.16 \text{ mm}$
 $D_3 = 2.82 \text{ mm}$
 $D_4 = 3.56 \text{ mm}$
 $t = 0.38 \text{ mm}$



³ Munson, B. R., Young, D. F., Okiishi, T. H., Fundamental Fluid Mechanics, 5th Edition, 2005.

⁴ Burattini, P. and Talamelli, A., "Acoustic control of a coaxial jet," J. of Turbulence, 8, 1-14, 2009.

Experimental Techniques: Image Analysis

- Proper Orthogonal Decomposition
 - For a set of high-speed images,

$$A = \sum_{k=1}^N q_k(t) \varphi_k(x)$$

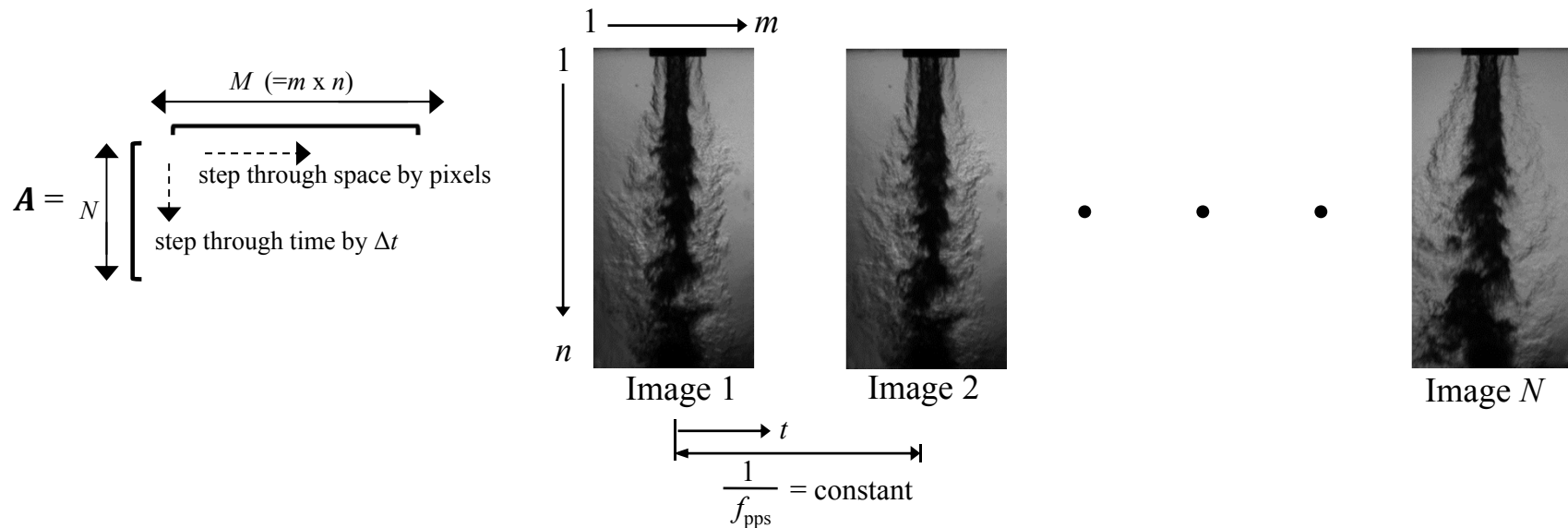
A : pixel intensity data matrix

q_k : vectors of temporal amplitude coefficients

φ_k : vectors of proper orthogonal modes

k : mode number

N : total number of modes



Experimental Techniques: Image Analysis

- Proper Orthogonal Decomposition
 - To identify periodic structures, subtract the average image

$$\tilde{A}_{ij} = A_{ij} - \frac{1}{N} \sum_i A_{ij} \quad i = 1 \dots N, \quad j = 1 \dots M$$

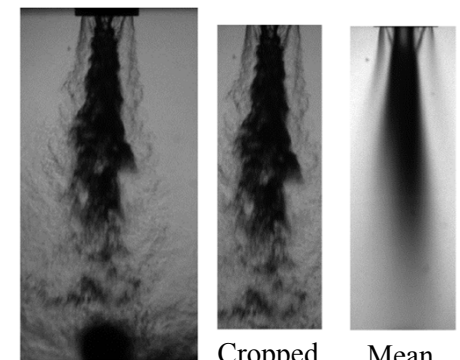
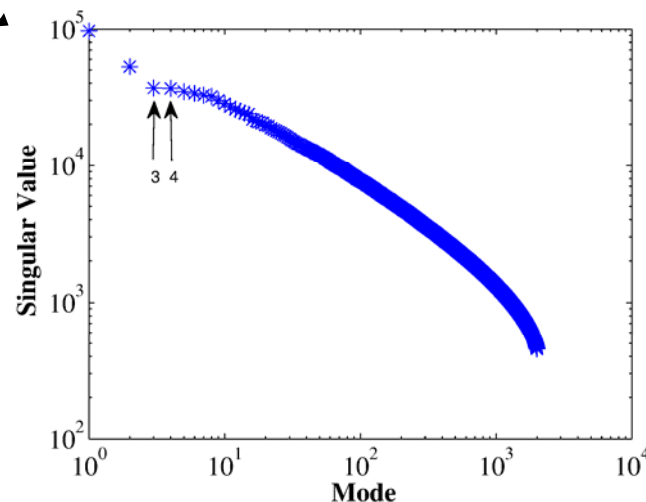
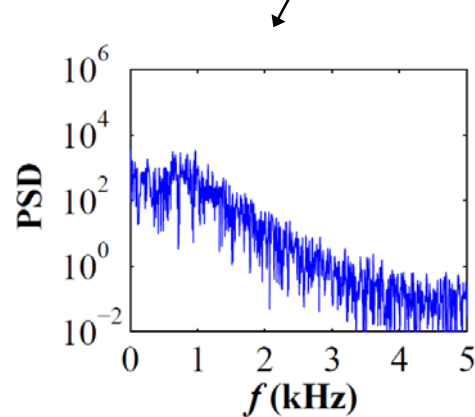
\tilde{A} : matrix of intensity fluctuations

- Singular Value Decomposition (SVD) of \tilde{A}

$$\tilde{A} = \tilde{U} \Sigma \tilde{V}^T$$

equivalent to $q_k(t)$

equivalent to $\varphi_k(x)$



Snapshot Image

Cropped

Mean

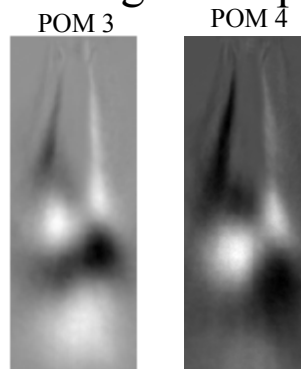
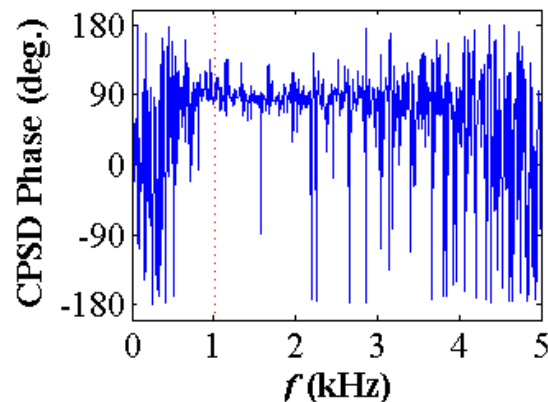
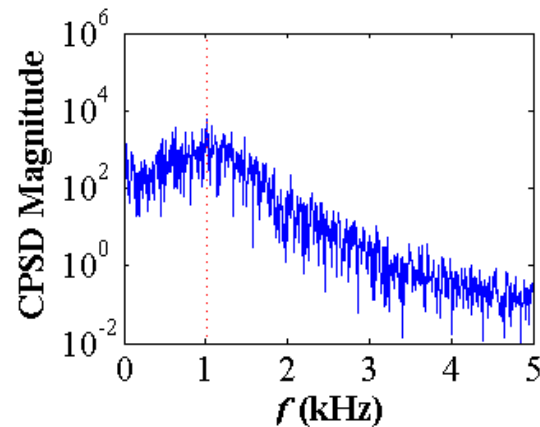
POM 3



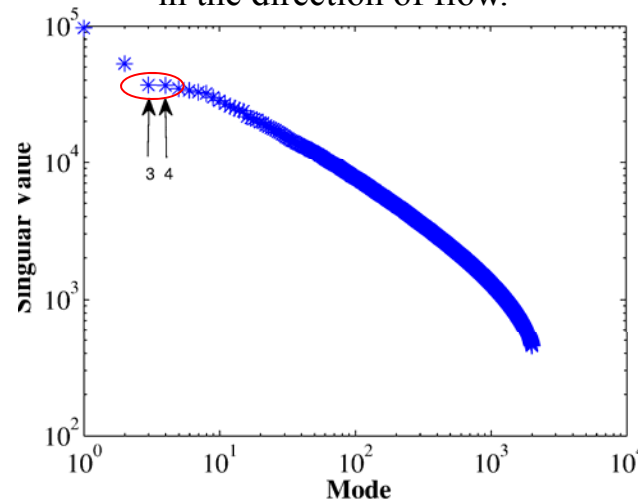
$p_c = 400$ psia
 $J = 6$

Experimental Techniques: Image Analysis

- Proper Orthogonal Decomposition
 - To identify traveling, coherent structures, a conjugate mode pair is identified as any two modes whose CPSD magnitude peaks near a phase of $\pm 90^\circ$.⁵



POMs 3 & 4 exhibit the same flow structure, shifted by 90° in the direction of flow.



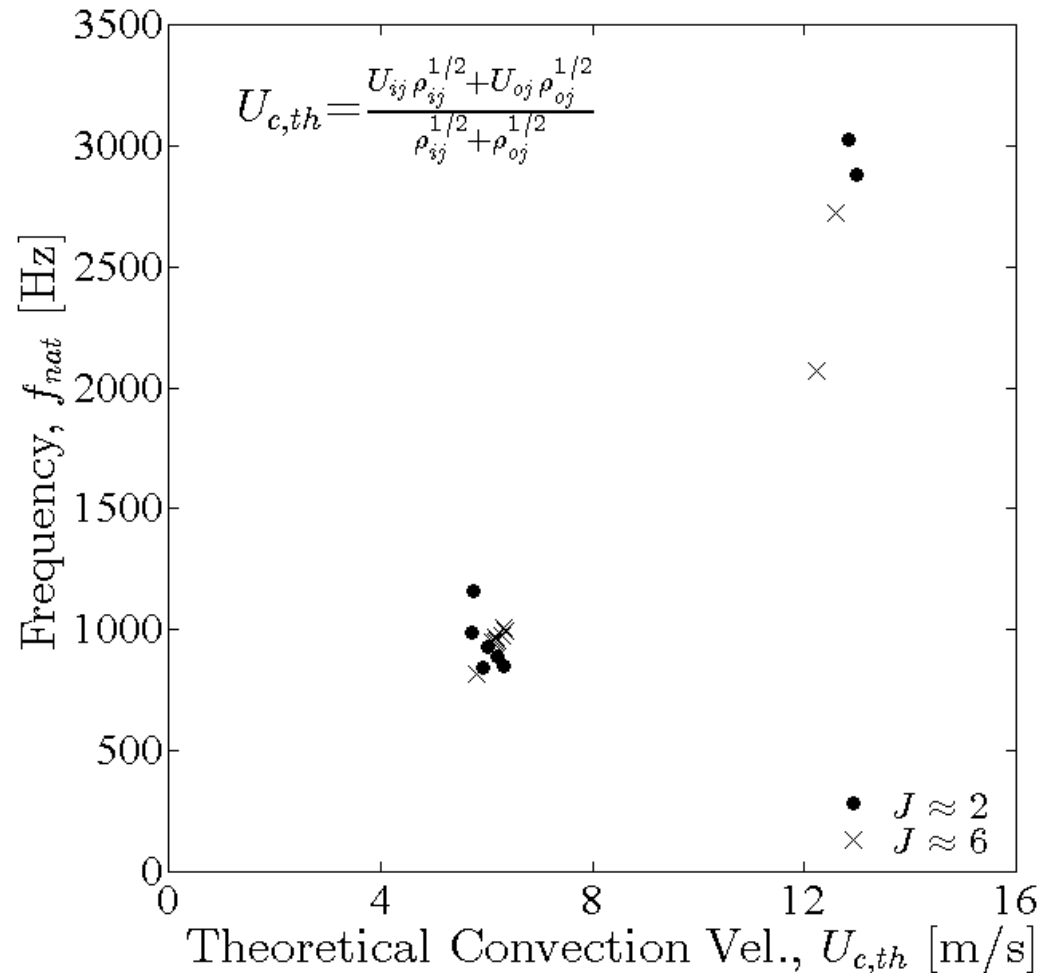
⇒ Proper orthogonal modes (POMs) 3 & 4 were found to be the **most energetic conjugate pair**.

⇒ The natural mode is represented by POMs 3 & 4.

⇒ The natural mode spans a band of frequencies rather than a single peak frequency.

Results: Unforced Coaxial Jets

- Natural Jet Characterization: f_{nat} values determined by POD



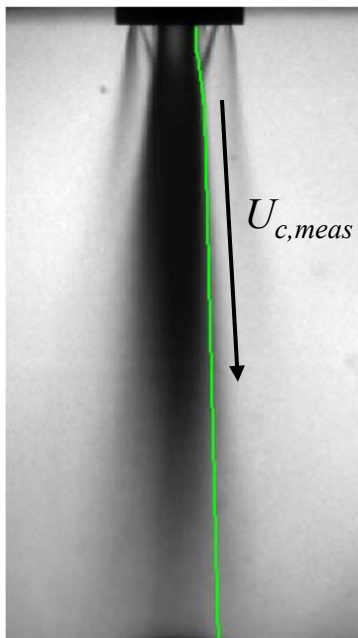
- Inner jet: LN_2
- Outer jet: GHe
- Two $U_{c,th}$ values are used
- Two J values are used
- $p_c = 400$ psia

$\Rightarrow f_{\text{nat}}$ scales with $U_{c,th}$
approximately

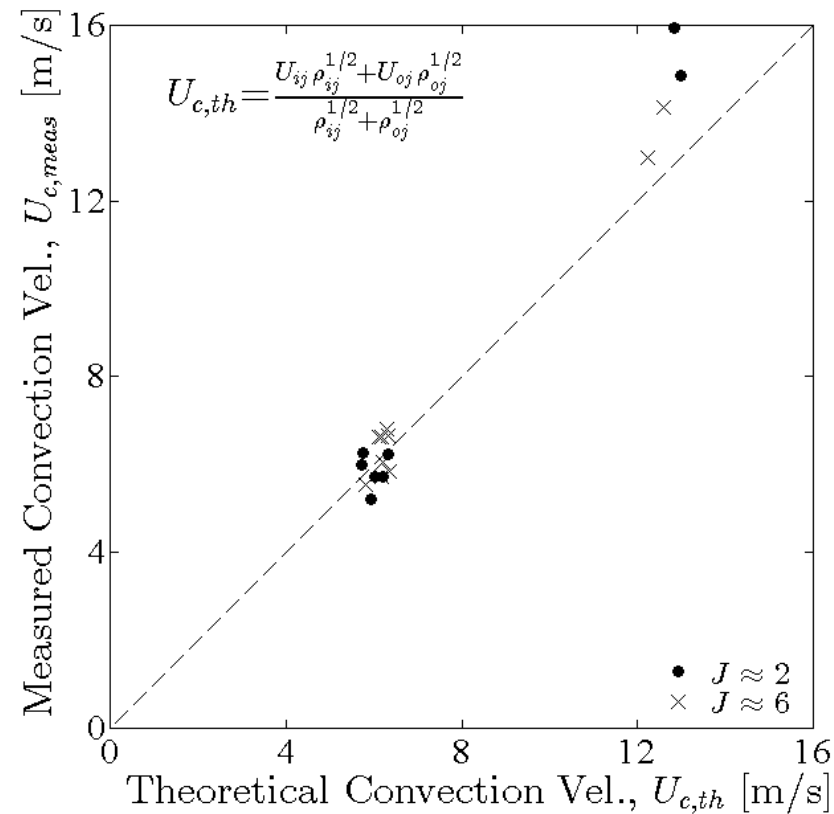
⁶ Dimotakis, P. E. 1986 "Two-Dimensional Shear-Layer Entrainment," *AIAA J.* 24, 1791-1796.

Results: Unforced Coaxial Jets

- High-speed images were used to **experimentally measure** the shear layer convection velocity



$$U_{c,meas} = \frac{\Delta s}{\Delta t}$$

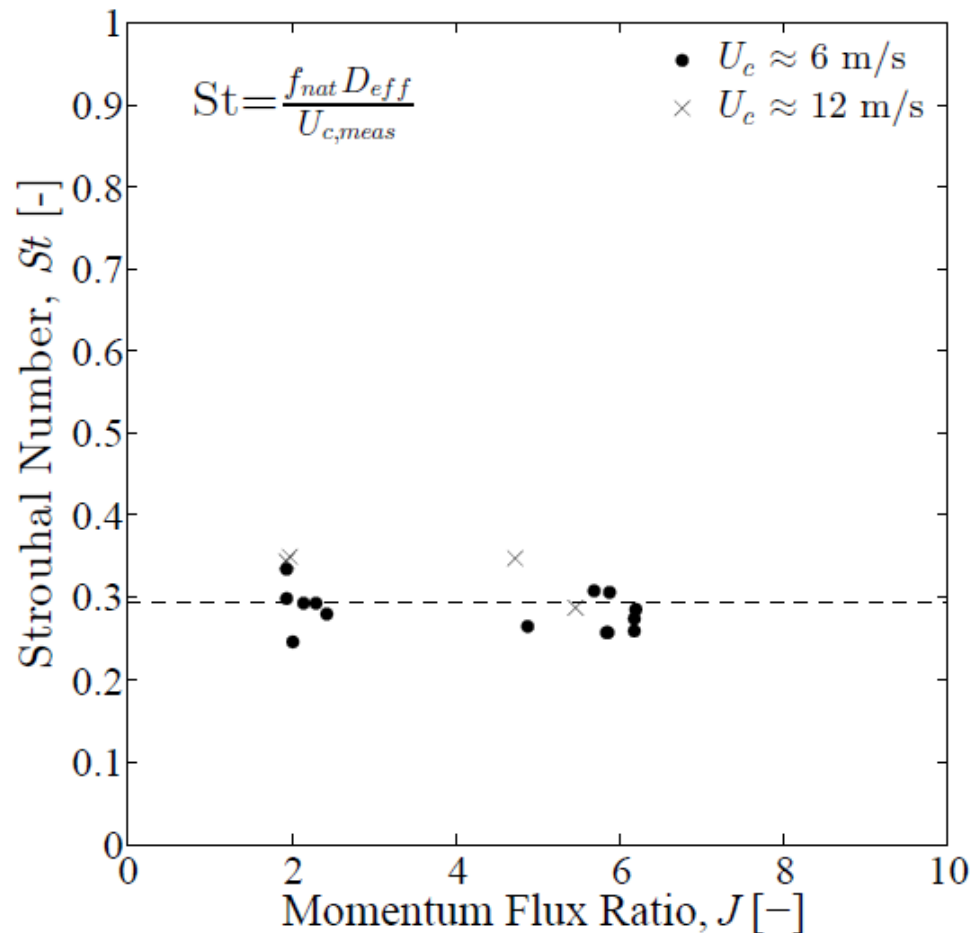


$\Rightarrow U_{c,th}$ accurately predicts convection velocities near 6 m/s

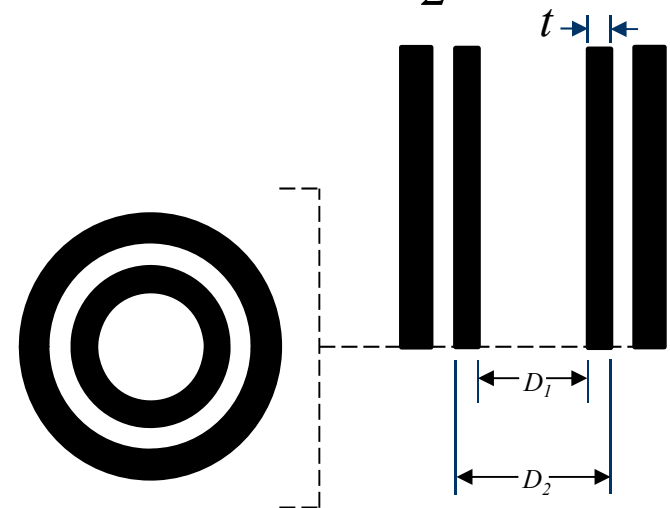
$\Rightarrow U_{c,th}$ under-predicts higher convection velocities

Results: Unforced Coaxial Jets

- Natural Jet Characterization: St scaling law produced by **experimental** shear layer convection velocities



- $$D_{eff} = \frac{D_1 + D_2}{2}$$

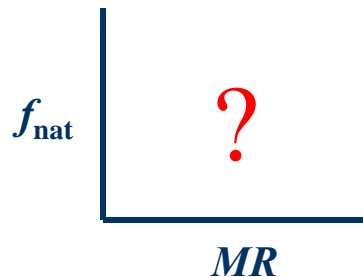


⇒ St values collapse around a mean of 0.29

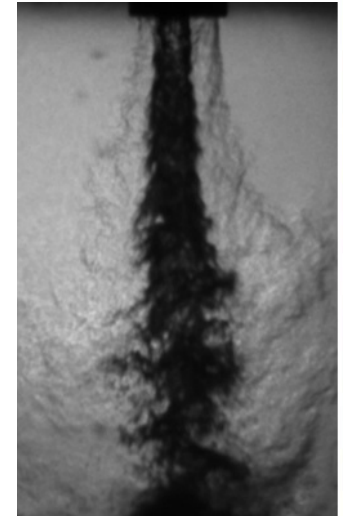
⇒ St is not dependent on J

Future Work

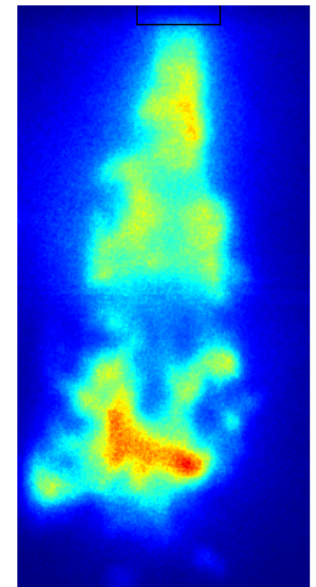
- Nonreactive Experiments
 - Explore susceptibility of jets to acoustic forcing with regard to:
 - Nondimensional Forcing Frequency $F = f_F/f_{\text{nat}}$
 - Acoustic pressure amplitude p'
 - Momentum Flux Ratio J
 - Injector Geometry $AR = A_{oj}/A_{ij}$
 - Explore the existence of convectively unstable and absolutely unstable coaxial jets
- Reactive Experiments
 - Characterize the spectral content of natural flame instabilities
 - Explore hydrogen and hydrocarbon fuels
 - Explore variations in injector geometry, including gas-centered swirl-coaxial injectors



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Non-reactive



Reactive

Summary

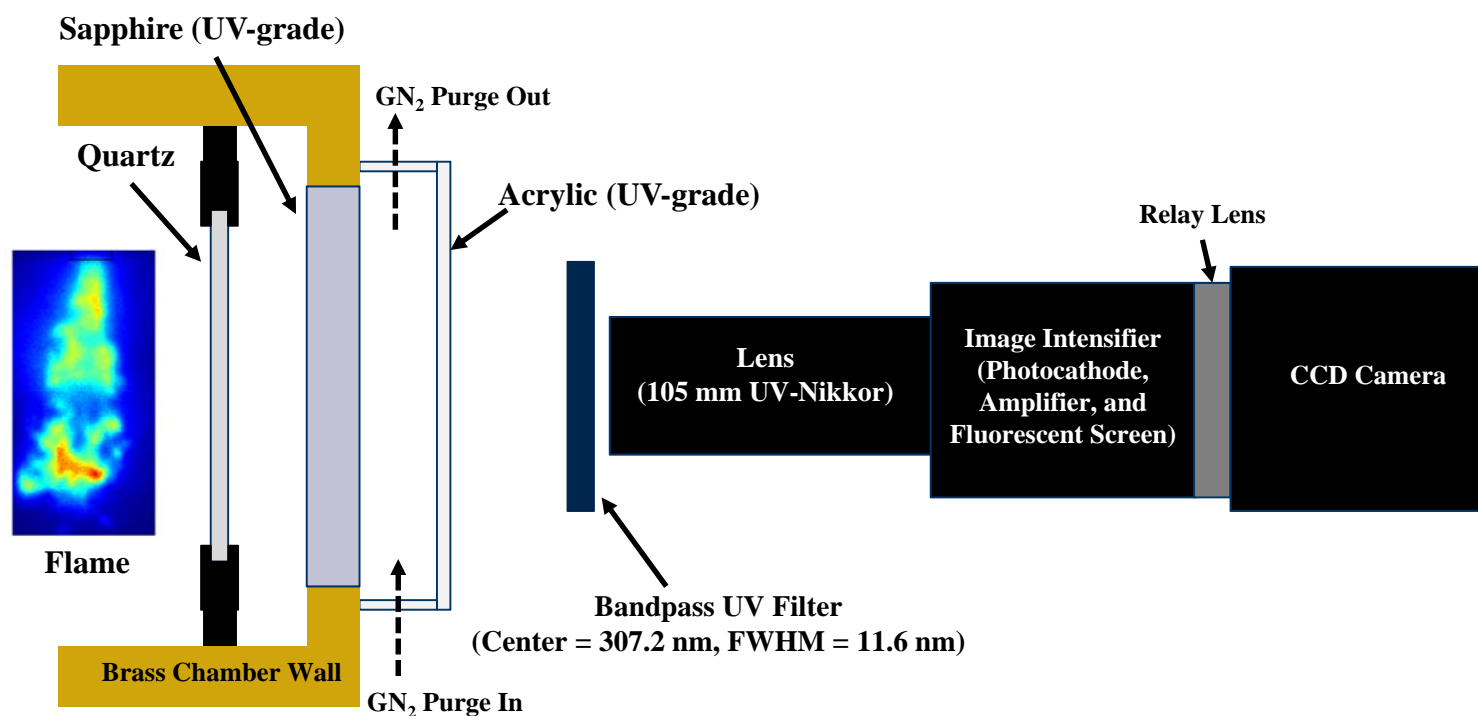
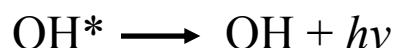
- A facility for combustion stability experiments was constructed to allow for chamber pressures up to 1500 psi and acoustic pressure amplitudes greater than 2% of the mean chamber pressure.
- Cryogenic heat exchangers were characterized as an effective technique to control the temperature of both oxygen and hydrogen at the injector.
- A shear-coaxial injector was designed and fabricated to produce fully developed turbulence at the exit for a wide range of Re.
- Preliminary, nonreactive results were analyzed using proper orthogonal decomposition of high-speed images to extract the dominant instability frequency of each flow condition.
- Future work will characterize the spectral behavior of reactive coaxial jets with regard to fuel type and injector geometry, and study the susceptibility of these flows to transverse acoustic forcing.

References

1. Richecoeur, F., "Experiments and numerical simulations of interactions between transverse acoustic modes and cryogenic flames," PhD thesis, Ecole Centrale Paris, November 2006.
2. Pierce, A. D., Acoustics: An Introduction to Its Physical Principles and Applications, 2nd Edition, 360-363, 1991.
3. Munson, B. R., Young, D. F., Okiishi, T. H., Fundamental Fluid Mechanics, 5th Edition, 2005.
4. Burattini, P. and Talamelli, A., "Acoustic control of a coaxial jet," J. of Turbulence, 8, 1-14, 2009.
5. Arienti, M. and Soteriou, M.C., "Time resolved proper orthogonal decomposition of liquid jet dynamics," *Phys of Fluids*, 21 112104, 2009.
6. Dimotakis, P. E. "Two-Dimensional Shear-Layer Entrainment," *AIAA* 24, 1791-1796, 1986.

Experimental Techniques: Facility

- OH* Chemiluminescence Imaging
 - Radiative de-excitation of hydroxyl radicals emits ultraviolet light at a wavelength of approximately 308 nm.

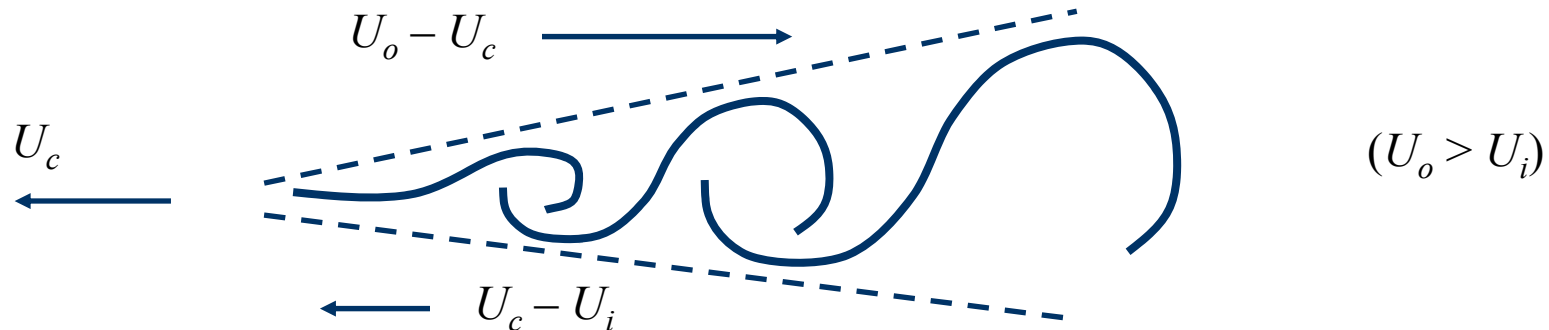


- Unwanted radiation from vibrational and rotational bands of H₂O products are not recorded.
- Nitrogen purge removes condensation from window surfaces.

Experimental Techniques: Test Matrix

Convective Shear Layer Velocity by Dimotakis [AIAA, 1986]⁵

- Vortex Frame of Reference



- Bernoulli's equation
 - A stagnation point must exist between vortices. Therefore, along a line through this point, dynamic pressures are approximately equal.

$$\rho_o (U_o - U_c)^2 \approx \rho_i (U_c - U_i)^2$$

$$U_c = \frac{U_o \rho_o^{1/2} + U_i \rho_i^{1/2}}{\rho_o^{1/2} + \rho_i^{1/2}}$$

$$St = \frac{f_{nat} D}{U_c}$$

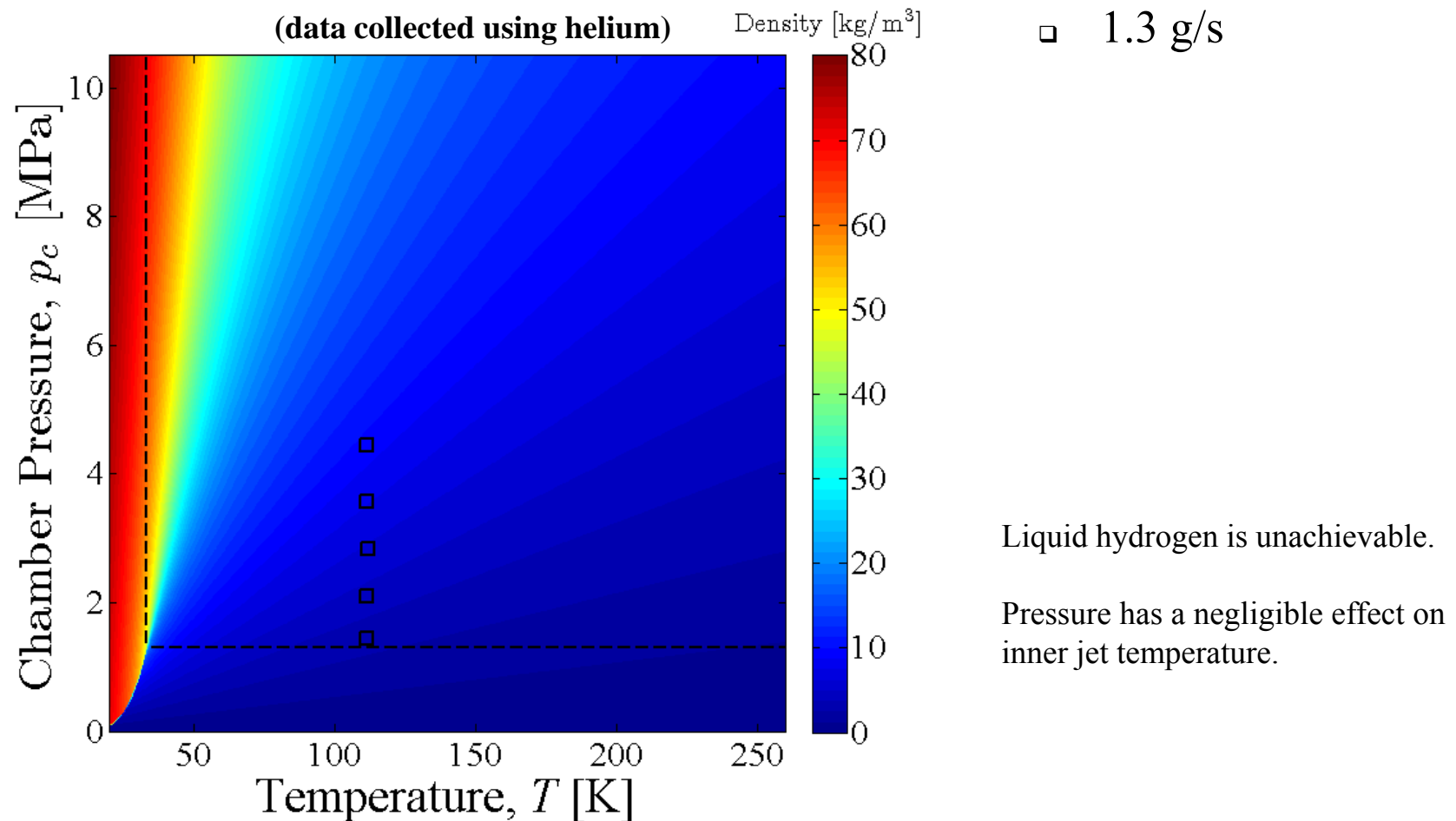


If St , D , U_c are held constant then f_{nat} may be constant.

⁵ Dimotakis, P. E. 1986 "Two-Dimensional Shear-Layer Entrainment," *AIAA J.* 24, 1791-1796.

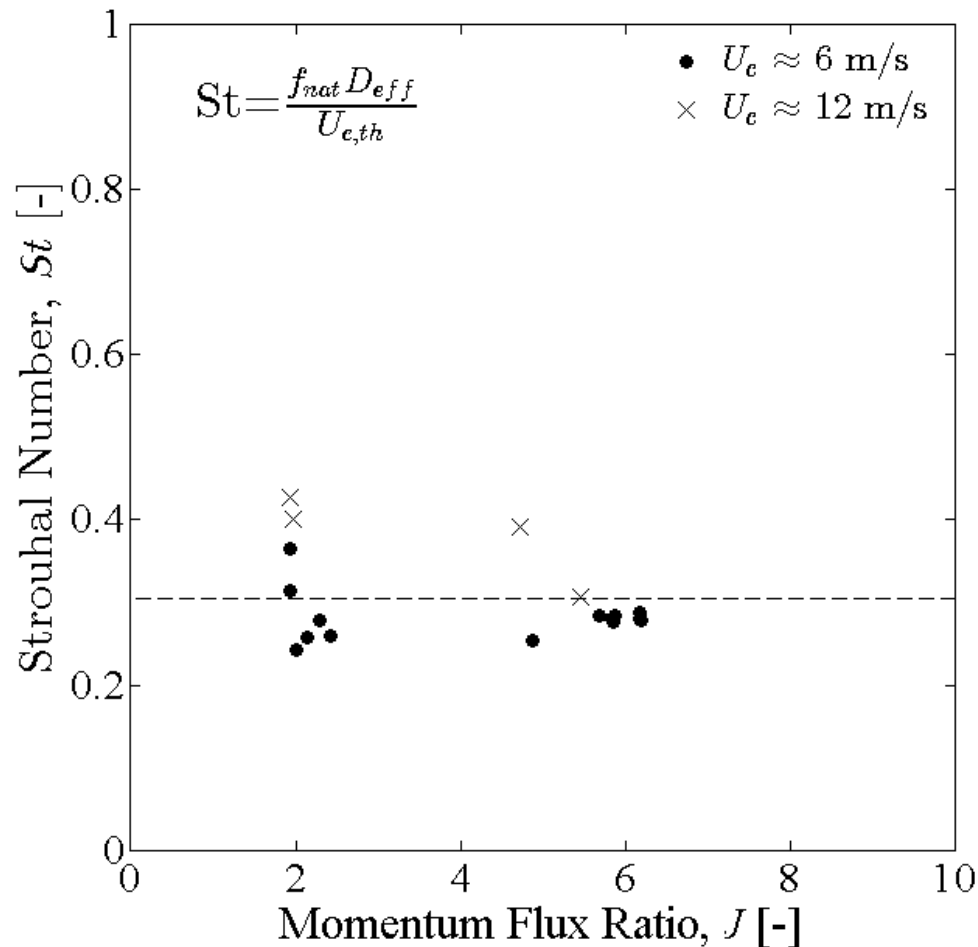
Experimental Techniques: Heat Exchangers

- Outer Jet Heat Exchanger

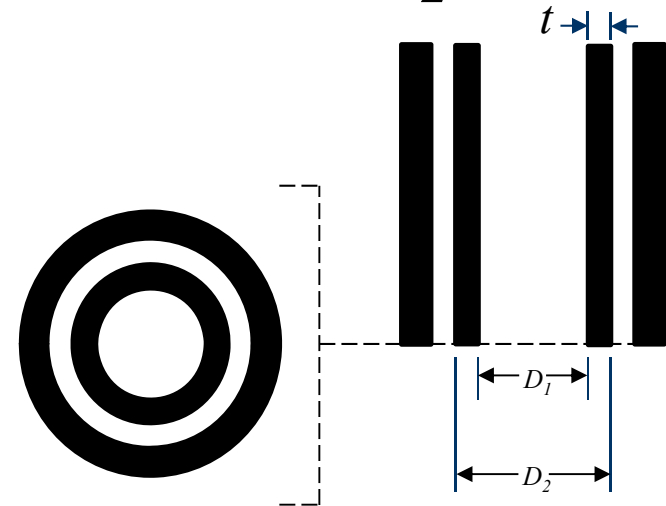


Results: Unforced Coaxial Jets

- Natural Jet Characterization: St scaling law produced by **theoretical** shear layer convection velocities



- $$D_{eff} = \frac{D_1 + D_2}{2}$$



- \Rightarrow St values have a mean of 0.30
- \Rightarrow St is not dependent on J
- \Rightarrow High shear layer convection velocities show more scatter